

**FÍSICA da MATÉRIA CONDENSADA**  
Mestrado em Engenharia Física Tecnológica  
Série 3b

1. Masses, Lagrangian, Hamiltonian and Legendre transformation.

a) Starting from the Lagrangian  $\mathcal{L}(\vec{x}, \vec{v}) = -mc^2\sqrt{1 - (\frac{\vec{v}}{c})^2} - V(\vec{x})$ , for a point particle, show that the equation of motion of the particle  $\frac{d\vec{p}}{dt} = -\nabla V$ , with  $\vec{p} = \frac{m\vec{v}}{\sqrt{1 - (\frac{\vec{v}}{c})^2}}$  can be written as  $m_{\perp}\vec{a}_{\perp} + m_{\parallel}\vec{a}_{\parallel} = -\nabla V$ , where  $(\vec{a}_{\perp})^i = \sum_j P_{\perp}^{ij} \frac{dv^j}{dt}$  and  $(\vec{a}_{\parallel})^i = \sum_j P_{\parallel}^{ij} \frac{dv^j}{dt}$ , with  $P_{\parallel}^{ij} = \frac{v^i v^j}{v^2}$  and  $P_{\perp} = I - P_{\parallel}$ . Determine the (transversal and longitudinal) masses  $m_{\perp}$  and  $m_{\parallel}$ .

b) Using for the energy the expression  $E = mc^2\sqrt{1 + (\frac{\vec{p}}{mc})^2} + V(\vec{x})$ , obtain  $\frac{\partial E}{\partial \vec{p}} = \vec{v}$  and  $\frac{\partial^2 E}{\partial p^i \partial p^j} = \frac{1}{m_{\perp}} P_{\perp}^{ij} + \frac{1}{m_{\parallel}} P_{\parallel}^{ij}$ , where  $P_{\parallel}^{ij} = \frac{p^i p^j}{p^2}$  and  $P_{\perp} = I - P_{\parallel}$ . What are  $m_{\perp}$  and  $m_{\parallel}$ , according to this definition? Verify that the masses and the projectors are the same as obtained above.

c) Explain this fact, using that the Hamiltonian is given by a Legendre transformation of the Lagrangian.

Hint:  $\frac{p}{mv} = \frac{1}{\sqrt{1 - (\frac{v}{c})^2}} = \sqrt{1 + (\frac{p}{mc})^2}$ .

2. Tight binding interaction and effective mass.

Consider a free electron system, in a d-dimensional lattice, with Hamiltonian given by

$$\mathcal{H} = \sum_{\alpha=1}^d \sum_{\vec{x}} a^{\dagger}(\vec{x}) t_{\alpha} a(\vec{x} + a\vec{e}_{\alpha}) + h.c. = \sum_{\vec{k}} a_{\vec{k}}^{\dagger} \epsilon(\vec{k}) a_{\vec{k}}$$

where  $a(\vec{x}) = \frac{1}{\sqrt{N}} \sum_{\vec{k}} e^{i\vec{k}\cdot\vec{x}} a_{\vec{k}}$  and  $a^{\dagger}(\vec{x}) = \frac{1}{\sqrt{N}} \sum_{\vec{k}} a_{\vec{k}}^{\dagger} e^{-i\vec{k}\cdot\vec{x}}$ , and  $t_{\alpha}$  is real and negative.

a) Obtain the band energy

$$\epsilon(\vec{k}) = \sum_{\alpha=1}^d 2t_{\alpha} \cos(k_{\alpha}a).$$

b) Define the effective masses  $m_{\alpha\beta}^{-1} = \frac{\partial^2 \epsilon(\vec{k})}{\partial k_{\alpha} \partial k_{\beta}}$ , for  $k_{\alpha} = 0$  and  $k_{\alpha} = \pm\pi/a$ . Interpret physically the results.

3. Density of states in a d-dimensional isotropic system.

a) In a d-dimensional system, in the thermodynamic limit, we can write

$$\frac{1}{V} \sum_{\vec{k}_n} \rightarrow \int \frac{d^d k}{(2\pi)^d} f(k) = \int_0^\infty S_d k^{d-1} f(k) \frac{d^d k}{(2\pi)^d}$$

where  $|k| = |\vec{k}|$  and  $S_d$  is the area of the sphere in a d-dimensional space.

Choosing  $f(k) = e^{-\frac{\vec{k}^2}{2}}$ , show that

$$S_d = \frac{2\pi^{d/2}}{\Gamma(\frac{d}{2})}$$

giving  $S_1 = 2$ ,  $S_2 = 2\pi$  and  $S_3 = 4\pi$ , for  $d = 1, 2, 3$ .

b) Considering a system of free particles with energy given by  $\epsilon(k) = \frac{(\hbar\vec{k})^2}{2m}$  (parabolic band), show that  $\mathcal{D}_d(\epsilon(k)) =$ , defined by

$$\int \frac{d^d k}{(2\pi)^d} f(\epsilon(k)) = \int_0^\infty \mathcal{D}_d(\epsilon) f(\epsilon) d\epsilon$$

is given by

$$\mathcal{D}_d(\epsilon) = \frac{S_d}{(2\pi)^d} k^{d-1} \frac{dk}{d\epsilon} = \frac{1}{2} \frac{S_d}{(2\pi)^d} k^{d-2} \frac{1}{\frac{d\epsilon}{dk^2}} = \frac{S_d}{(2\pi)^d} \frac{m}{\hbar^2} \left( \frac{2m\epsilon}{\hbar^2} \right)^{\frac{d-2}{2}}$$

c) Plot  $\mathcal{D}_d(\epsilon)$ , as a function of  $\epsilon$ , for  $d = 1, 2, 3$ .

4. Quantum gas in two dimensions.

Consider a free Bosonic or Fermionic system, in 2 dimensions, in thermal equilibrium at temperature  $T$ .

a) Write down the equations for  $N = \langle \hat{N} \rangle$  and  $E = \langle \mathcal{H} \rangle$ .

b) Solve the equation  $N = \langle \hat{N} \rangle$  for  $\mu$ .

c) Show that the specific heat of the Bosonic and Fermionic systems is the same.

References:

Quantum Statistics of Ideal Gases in Two Dimensions, R. M. May, Phys. Rev. **153**, A1515 (1964).

Specific Heat of Two-Dimensional Ideal quantum Gases, V. V. Ul'yanov and S. S. Sokolov, Soviet Physics Journal **18**, 138-139 (1975).

5. Itinerant magnetism at  $T = 0$ .

Consider a system with a fixed number  $N$  of electrons, described by the Hubbard model with Hamiltonian given by

$$\mathcal{H} = \sum_{ij} \sum_{\sigma} c_{i\sigma}^{\dagger} t_{ij} c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow},$$

without external magnetic field, at temperature  $T = 0$ .

a) Use the product of two Slater determinants, one for each value of the spin, or, equivalently, modify the usual Fermi sea, according to

$$|\Phi_{Stoner}\rangle = \prod_{k < k_{F\uparrow}} c_{k\uparrow}^{\dagger} \prod_{k < k_{F\downarrow}} c_{k\downarrow}^{\dagger} |0\rangle$$

for the ferromagnetic state with a uniform magnetization  $M = \frac{1}{2}(N_{\uparrow} - N_{\downarrow})$ , minimize the energy  $E = \langle \mathcal{H} \rangle$ , subject to the constraint  $N = N_{\uparrow} + N_{\downarrow}$ , introducing a Lagrange multiplier  $\mu$ , leading to the equations

$$\begin{aligned} \epsilon(k_{F\uparrow}) - \mu + U n_{\downarrow} &= 0 \\ \epsilon(k_{F\downarrow}) - \mu + U n_{\uparrow} &= 0 \end{aligned}$$

where  $n_{\uparrow} = \frac{N_{\uparrow}}{V}$ ,  $n_{\downarrow} = \frac{N_{\downarrow}}{V}$ , giving

$$\epsilon(k_{F\uparrow}) + U n_{\downarrow} = \epsilon(k_{F\downarrow}) + U n_{\uparrow}$$

together with

$$N = N_{\uparrow} + N_{\downarrow}.$$

Interpret physically.

b) Considering free electrons with energy given by  $\epsilon(k) = \frac{(\hbar k)^2}{2m}$  show that

$$\frac{2\mathcal{D}\frac{E}{V}}{n^2} = \frac{9}{20} \left[ \left(1 + \frac{2m}{n}\right)^{5/3} + \left(1 - \frac{2m}{n}\right)^{5/3} \right] + \frac{1}{2} \mathcal{D}U \left(1 - \left(\frac{2m}{n}\right)^2\right)$$

where  $n = \frac{N}{V}$ ,  $m = \frac{M}{V}$  and  $\mathcal{D}$  is the density of states at  $k_F$ , in the absence of interactions i.e. for  $U = 0$ .

Minimizing the energy, verify that there is a trivial solution  $m = 0$  and, using the nontrivial solution, find the Stoner condition  $\mathcal{D}U > 1$  for the existence of the phase transition at  $T = 0$ .

c) Show that the magnetic phase is energetically more stable, when the Stoner criterion is satisfied.