

## Lecture 5

# Transport in low-dimensional systems: size quantization effects

- Two-dimensional electron gas
- Semiconductor quantum wells
- Quantum wires
- Quantum dots
- Spin-orbit interaction in low-dimensional systems

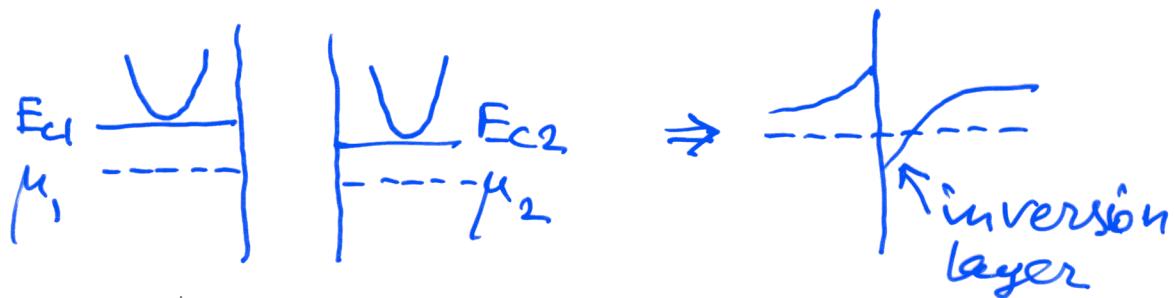
# Transport in low-dimensional systems

1

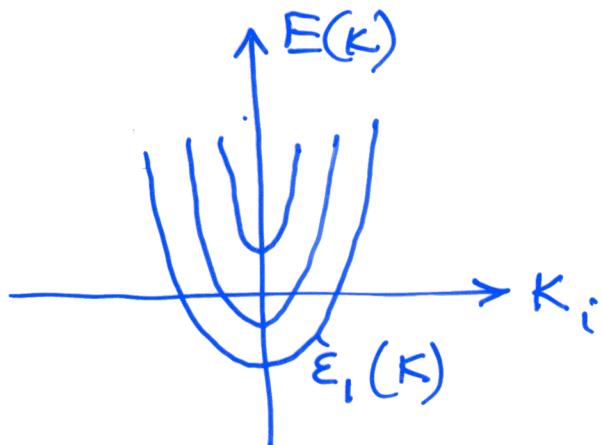
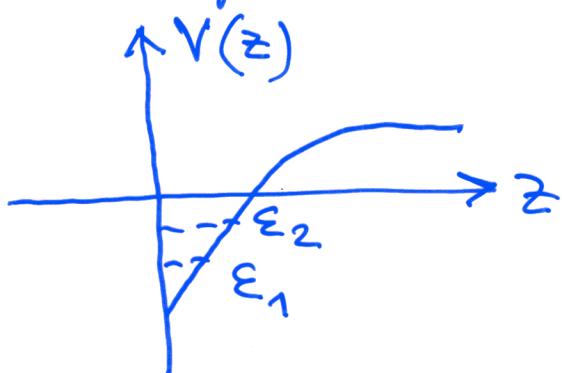
Two-dimensional electron gas (2DEG)

- Si-inversion layers (MOSFET technology)
- semiconductor interfaces (InGaAs-InAlAs)
- semiconductor quantum wells (QW)
- electrons on liquid helium

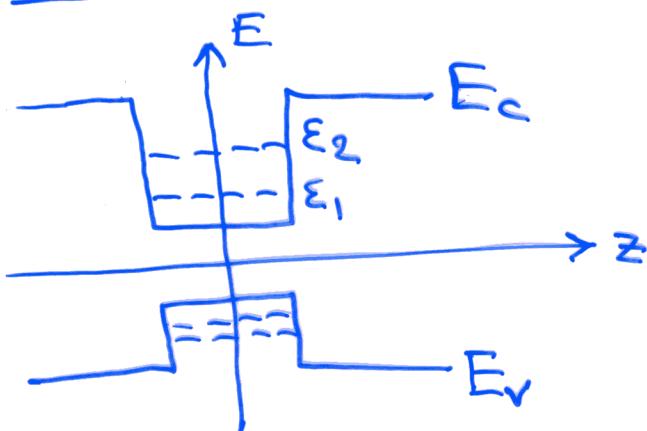
Contact of two semiconductors:



Potential profile



## Semiconductor QWs



Doping and scattering are separated (L. Esaki)

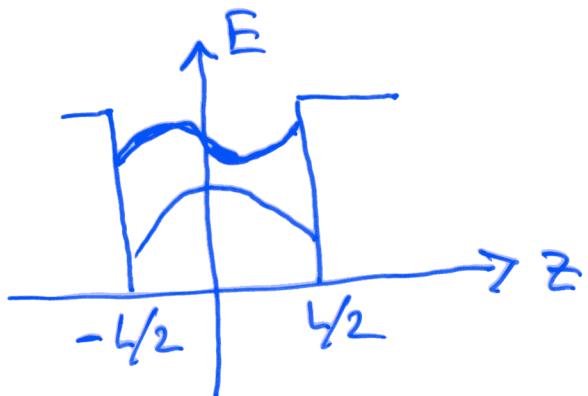
Si/SiGe, GaAlAs/GaAs  
PbTe/PbSnTe, etc.

rectangular model potential

Deep QW:

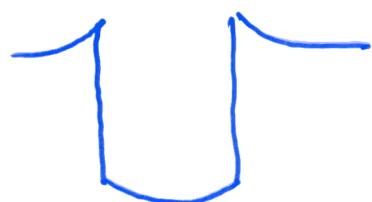
$$kL = n\pi$$

$$E_n = \frac{\hbar^2 n^2 \pi^2}{2mL^2}$$



$\langle \sim \rangle$  - condition of size-quantization

Real profile:  
(example)



Transport in 2D system: usual classical methods including Drude-Lorentz formula, Kubo formalism

(3)

## Problems:

- self-consistent potential profile
- scattering from impurities (screened potential)
- complex energy spectrum of semiconds.

## Estimations:

metals:  $\lambda \sim \lambda_0$ 

$$n = 2 \frac{4\pi k_F^3}{3(2\pi)^3}$$

Semiconductors:

$$\text{Let } n = 10^{15} \text{ cm}^{-3}$$

$$k = (3\pi^2 n)^{1/3} \approx 3 \cdot 10^5 \text{ cm}^{-1}$$

$$\lambda \approx 10^{-5} \text{ cm} = 100 \text{ nm}$$

Nanotechnology  $L \sim 10 \text{ nm}$ 

Low-dimensionality of a different type:

 $L \sim l$  (mean free path)

In Semiconductors:

$$l = v\tau \approx 10^7 \frac{\text{cm}}{\text{s}} \cdot 10^{-13} \text{ s} = 10^{-6} \text{ cm} = 10 \text{ nm}$$

$(h \approx 10^{18})$

In metals:

$$l = 10^8 \frac{\text{cm}}{\text{s}} \cdot 10^{-13} = 10^{-5} \text{ cm} = 100 \text{ nm}$$

## Effect of magnetic field in QW.

- parallel
- perpendicular

In perpendicular field:

- classically calculated Hall conductivity and magnetoresistance
- strong field: RHE

In parallel field:

- if  $l_K \leq$  magn. field is strong
- otherwise: weak corrections

$$l_K = \sqrt{\frac{hc}{eH}}$$

In perpendicular field:

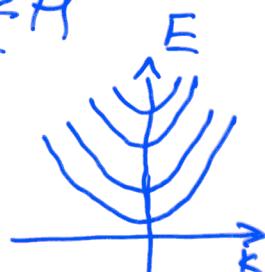
magn. field effect is separated

$$\epsilon_n(k) = \frac{\hbar^2 k^2}{2m} + \hbar \omega_K \left(n + \frac{1}{2}\right) \pm \frac{g\mu_B}{2} H$$

$$\omega_K = \frac{eH}{mc}$$

More complicated quasi 2D systems

- multilayers
- hybrid structures
- layered crystals (GaSe, InSe)



# Energy spectrum of QW in parallel magnetic field

(5)

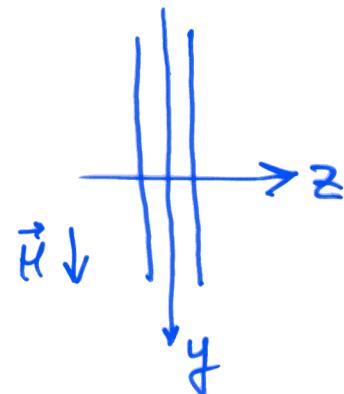
Schödinger eq:

$$\left[ -\frac{\hbar^2}{2m} \left( \nabla - \frac{ie\vec{A}}{\hbar c} \right)^2 + \frac{m\omega_0^2 z^2}{2} \right] \psi = \varepsilon \psi$$

parabolic QW

Let  $\vec{H} \parallel y$ :

$$\vec{A} = (H_z, 0, 0) \quad \begin{array}{l} \text{Landau} \\ \text{gauge} \end{array}$$



$$\left[ -\frac{\hbar^2}{2m} \left( \partial_x - \frac{ieH_z}{\hbar c} \right)^2 + \frac{\hbar^2}{2m} (\partial_y^2 + \partial_z^2) + \frac{m\omega_0^2 z^2}{2} - \varepsilon \right] \psi = 0$$

$\omega_0$  is classical frequency:

$$m\ddot{z} = -\omega_0^2 z m$$

$$z \sim e^{-i\omega_0 t}$$

$$V(z) = \frac{m\omega_0^2 z^2}{2}$$

$$\psi \sim e^{ik_x x + ik_y y}$$

$$\left[ \frac{\hbar^2 (k_x^2 + k_y^2)}{2m} - \frac{\hbar K_x eH_z}{mc} + \underbrace{\frac{e^2 H^2 z^2}{2mc^2}}_{\text{Harmonic potential}} - \frac{\hbar^2}{2m} \frac{d^2}{dz^2} + \underbrace{\frac{m\omega_0^2 z^2}{2}}_{\text{Harmonic potential}} - \varepsilon \right] \psi = 0$$

$$\left[ \frac{\frac{\hbar^2}{2m}(k_x^2 + k_y^2)}{2m} + \frac{m\omega^2}{2} \left( z - \frac{\omega_B \hbar k_x}{\omega^2 m} \right)^2 - \frac{\hbar^2 \omega_B^2 k_x^2}{2m \omega^2} - \frac{\hbar^2}{2m} \frac{d^2}{dz^2} - \varepsilon \right] \psi = 0$$

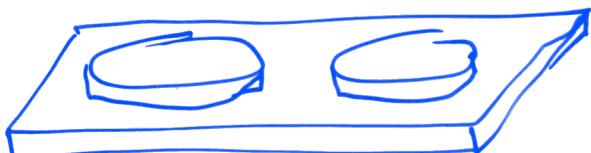
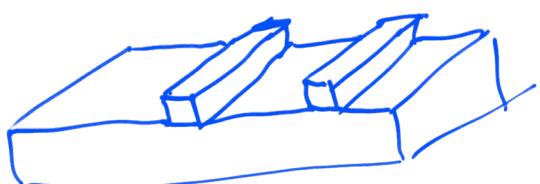
$$\left[ \frac{\frac{\hbar^2 K_x^2}{2m}}{2m} \frac{\omega_0^2}{\omega^2} + \frac{\frac{\hbar^2 K_y^2}{2m}}{2m} + \frac{m\omega^2}{2} (z - z_0)^2 - \frac{\hbar^2}{2m} \frac{d^2}{dz^2} - \varepsilon \right] \psi = 0$$

$\omega^2 = \omega_B^2 + \omega_0^2$

$$E(k) = \frac{\frac{\hbar^2 K_y^2}{2m}}{2m} + \frac{\frac{\hbar^2 K_x^2}{2m}}{2m} \frac{\omega_0^2}{\omega^2} + \hbar\omega(n + \frac{1}{2})$$

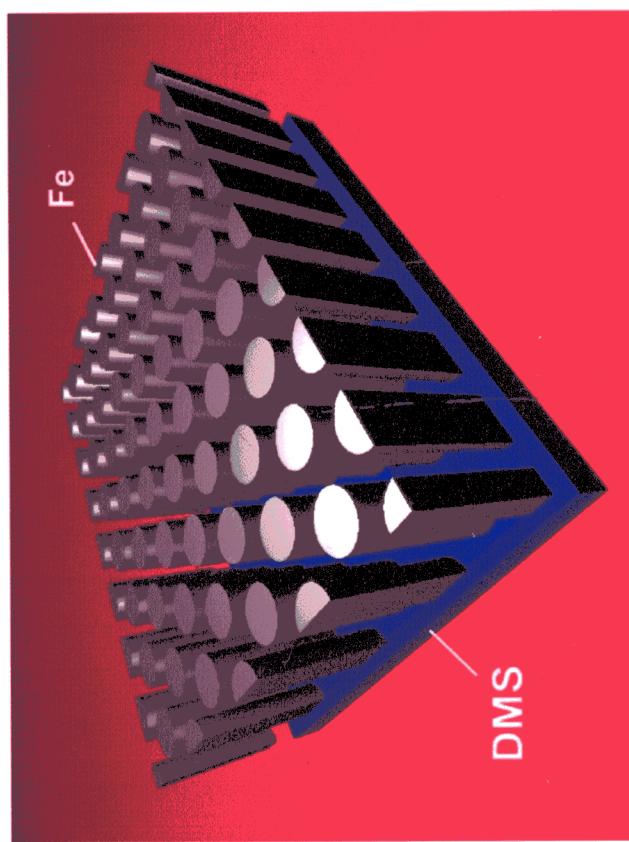
Other samples of low-dimensional objects:

- nanowires (semiconductor nanowires, carbon nanotubes)
- quantum dots



- nanowire lattice
- arrays of quantum dots etc.

## Magnetic nanowires



K. Nielsch et al. APL (2001)

# 2DEG with Rashba SO interaction

(7)

$$H = \frac{\hbar^2 k^2}{2m} + \alpha (\sigma_x K_y - \sigma_y K_x) =$$

$$= \begin{pmatrix} \epsilon_k & i\alpha K_- \\ -i\alpha K_+ & \epsilon_k \end{pmatrix}$$

$$\epsilon_k = \frac{\hbar^2 k^2}{2m}$$

(electrons in x-y plane)

$$K_{\pm} = K_x \pm iK_y$$

Schrödinger eq:

$$\begin{pmatrix} \epsilon_k - E & i\alpha K_- \\ -i\alpha K_+ & \epsilon_k - E \end{pmatrix} \begin{pmatrix} \psi \\ \chi \end{pmatrix} = 0$$

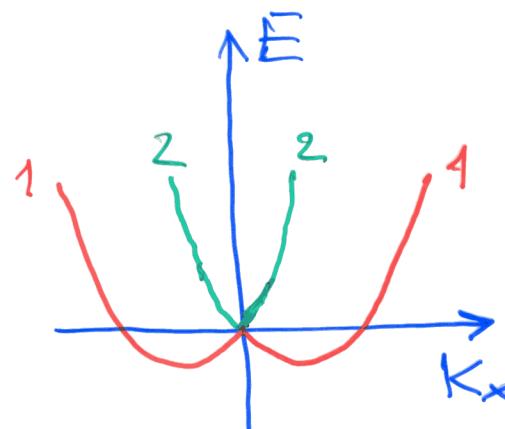
$$\sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$\Psi = \begin{pmatrix} \psi \\ \chi \end{pmatrix}$$

$$\begin{cases} (\epsilon_k - E)\psi + i\alpha K_- \chi = 0 \\ -i\alpha K_+ \psi + (\epsilon_k - E)\chi = 0 \end{cases}$$

$$\det = (\epsilon_k - E)^2 + \alpha^2 K^2 = 0$$

$$E_{1,2}(k) = \epsilon_k \mp \alpha K$$



Eigenfunctions:

$$1) \text{ For } E_1(k) = \epsilon_k - \alpha K$$

$$\chi = \frac{i\alpha K_+}{\epsilon_k - E_1} \quad \psi = \frac{iK_+}{K} \psi$$

$$\Psi_1 = N \begin{pmatrix} K \\ iK+ \end{pmatrix} \rightarrow \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ \frac{iK+}{K} \end{pmatrix}$$

$$\text{If } \vec{K} \parallel x : \quad \Psi_1 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ i \end{pmatrix}$$

Eigenfunction of  $\sigma_y$ :

$$\begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 1 \\ i \end{pmatrix} = \begin{pmatrix} 1 \\ i \end{pmatrix}$$

2)  $E_2(k) = \varepsilon_k + \Delta k$

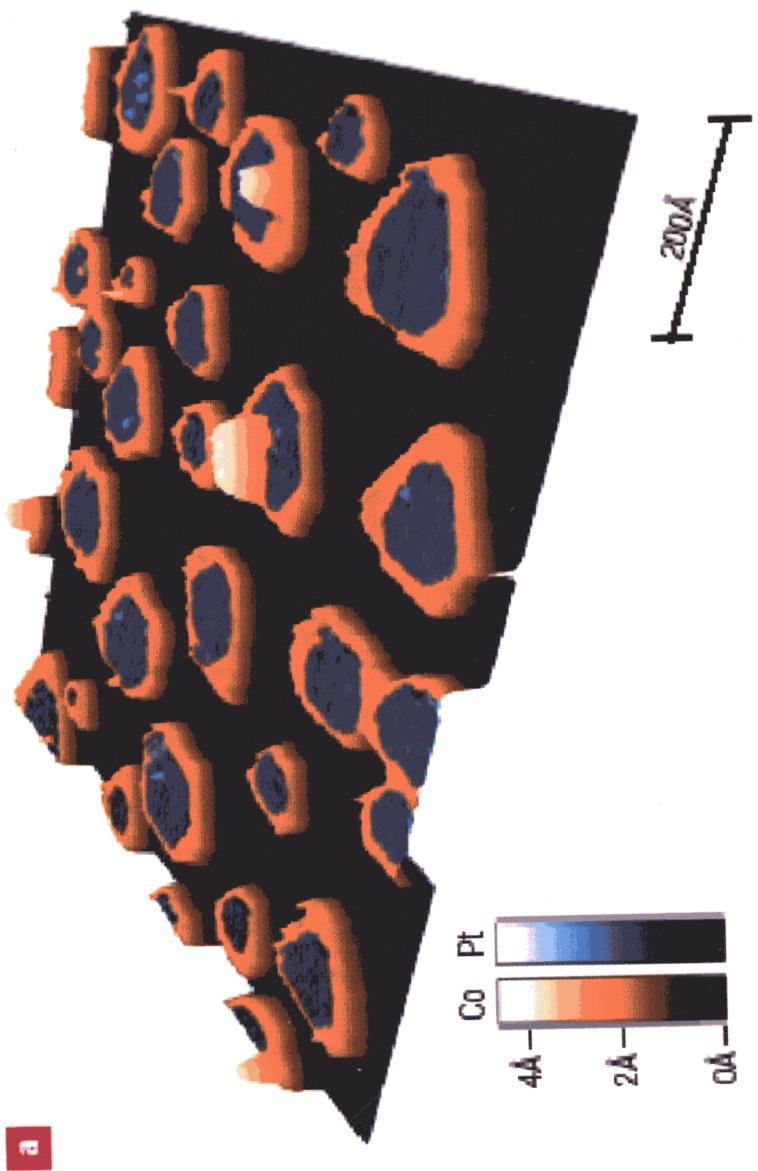
$$\varphi = \frac{i\Delta k_-}{E_2 - \varepsilon_k} = \frac{iK_-}{K} \chi$$

$$\Psi_2 = \frac{1}{\sqrt{2}} \begin{pmatrix} \frac{iK_-}{K} \\ 1 \end{pmatrix}$$

$$\text{If } \vec{K} \parallel x : \quad \Psi_2 = \frac{1}{\sqrt{2}} \begin{pmatrix} i \\ 1 \end{pmatrix}$$

It is also eigenfunction of  $\sigma_y$ :

$$\begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} i \\ 1 \end{pmatrix} = - \begin{pmatrix} i \\ 1 \end{pmatrix}$$



Tailoring bimetallic islands  
S. Rusponi et al. Nature (2003)

