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Non-ordinary nature Regge Trajectory of the $f_0(500)$ or σ -meson

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In collaboration with:

J. Nebreda, A. Szczepaniak and T. Londergan, Phys. Lett. B 729 (2014) 9–14
and J.A. Carrasco in preparation.

EEF70:

Workshop on Unquenched hadron spectroscopy: Non-perturbative Models and Methods of QCD versus experiment
Coimbra 1-5 September 2014

Motivation

- Interest in identification of non-ordinary Quark Model states (non $q\bar{q}$?)
- “Easy” if quantum numbers are not $q\bar{q}$ -> Exotics!
- Not so easy for cryptoexotics like light scalars.

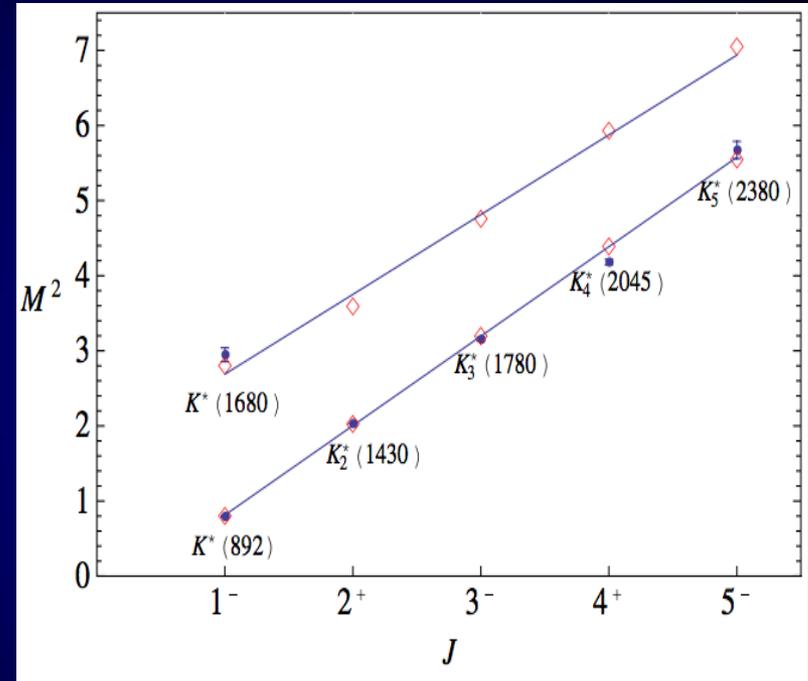
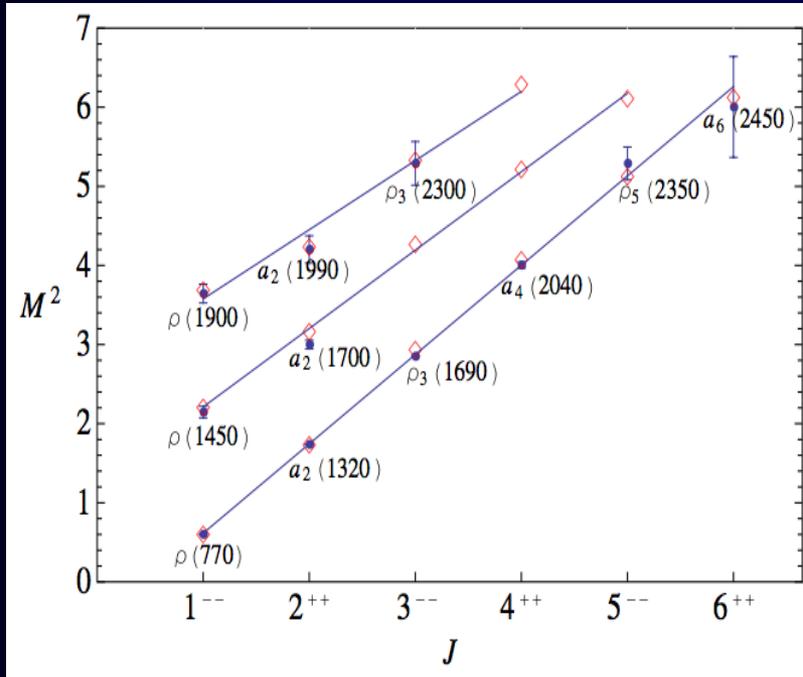
Particularly the $f_0(500)$ or σ -meson nature has been debated for over 50 years despite being very relevant for NN attraction, chiral symmetry breaking, Glueball search, lots of decays, etc...

- Hard to tell what a non-ordinary meson resonance is: tetraquark, molecule, glueball...
- Classification in terms of SU(3) multiplets complicated by **mixing**.

But “ordinary” $q\bar{q}$ mesons also follow another classification

Introduction: Regge trajectories

Particles with same quantum numbers and signature ($\tau=(-1)^J$) can be classified in **linear trajectories of (mass)² vs. (spin)** with a “universal” slope of $\sim 1 \text{ GeV}^{-2}$



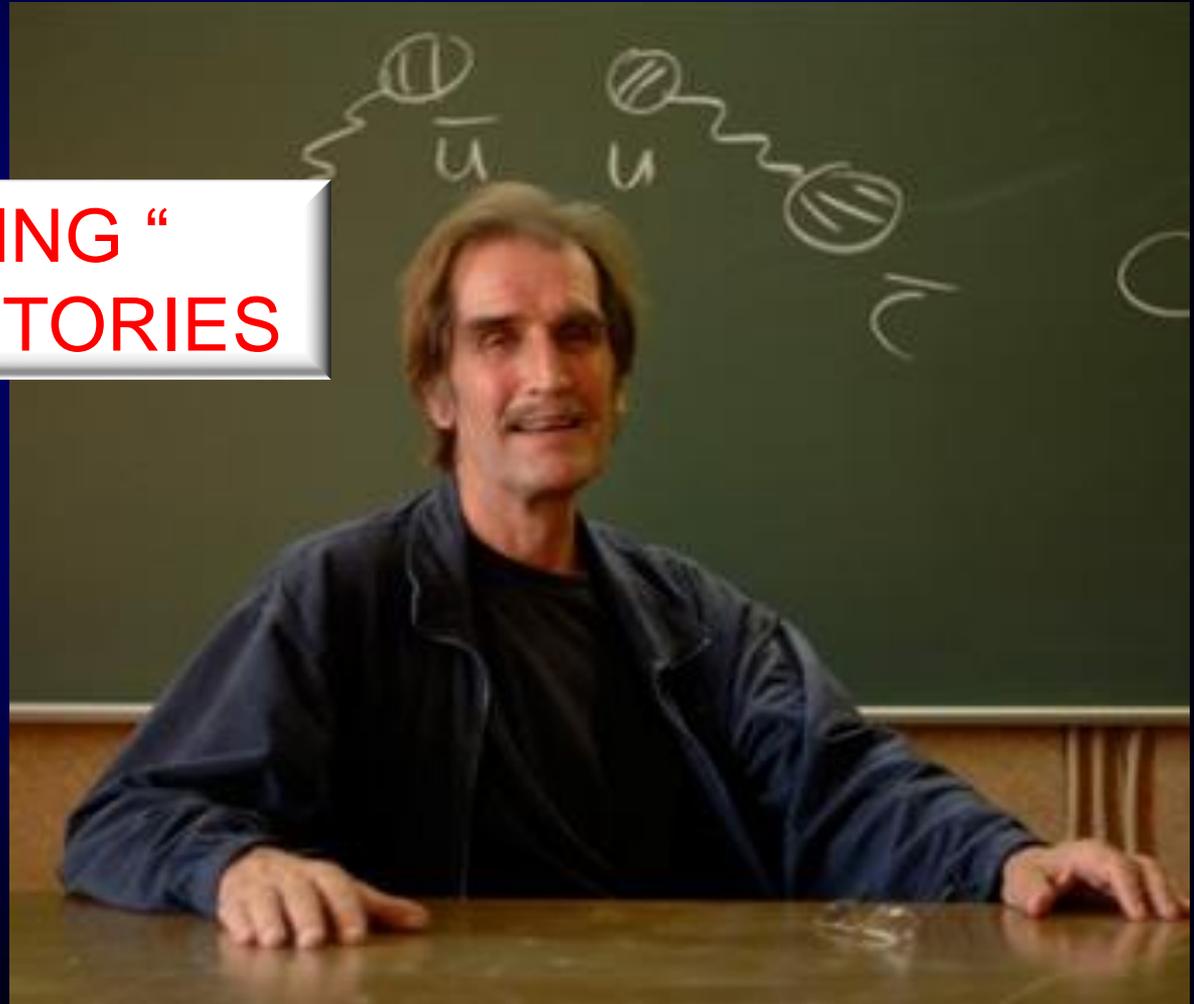
Anisovich-Anisovich-Sarantsev-Phys.Rev.D62.051502-4

Warning..., resonances have a width, and this variable here is... “ only M^2 ”
...some authors use width as uncertainty

In this work we aimed at including the width properly

Let us then dedicate this talk to Eef
and, for this talk, play
a little with the title
and change it to...

**“UNQUENCHING “
REGGE TRAJECTORIES**



But, hey!! The name is only
valid while in Coimbra



Introduction: Regge trajectories

Particles on each trajectory are somehow related by similar dynamics

Quark-antiquark states are well accommodated in these trajectories

A relativistic spinning rod of constant mass/length also has a linear spectrum

A color flux tube between a quark and an antiquark whose energy grows linearly with quark separation (confinement) could mimic this rod and is a crude model for Regge-linear trajectories (see Greensite's textbook) -> string theory

**Thus “ordinary” mesons, usually identified as $q\bar{q}$ states,
fit within
linear Regge trajectories**

But if other resonances have different nature...

.... they do not have to fit well in this scheme

Introduction: Regge trajectories

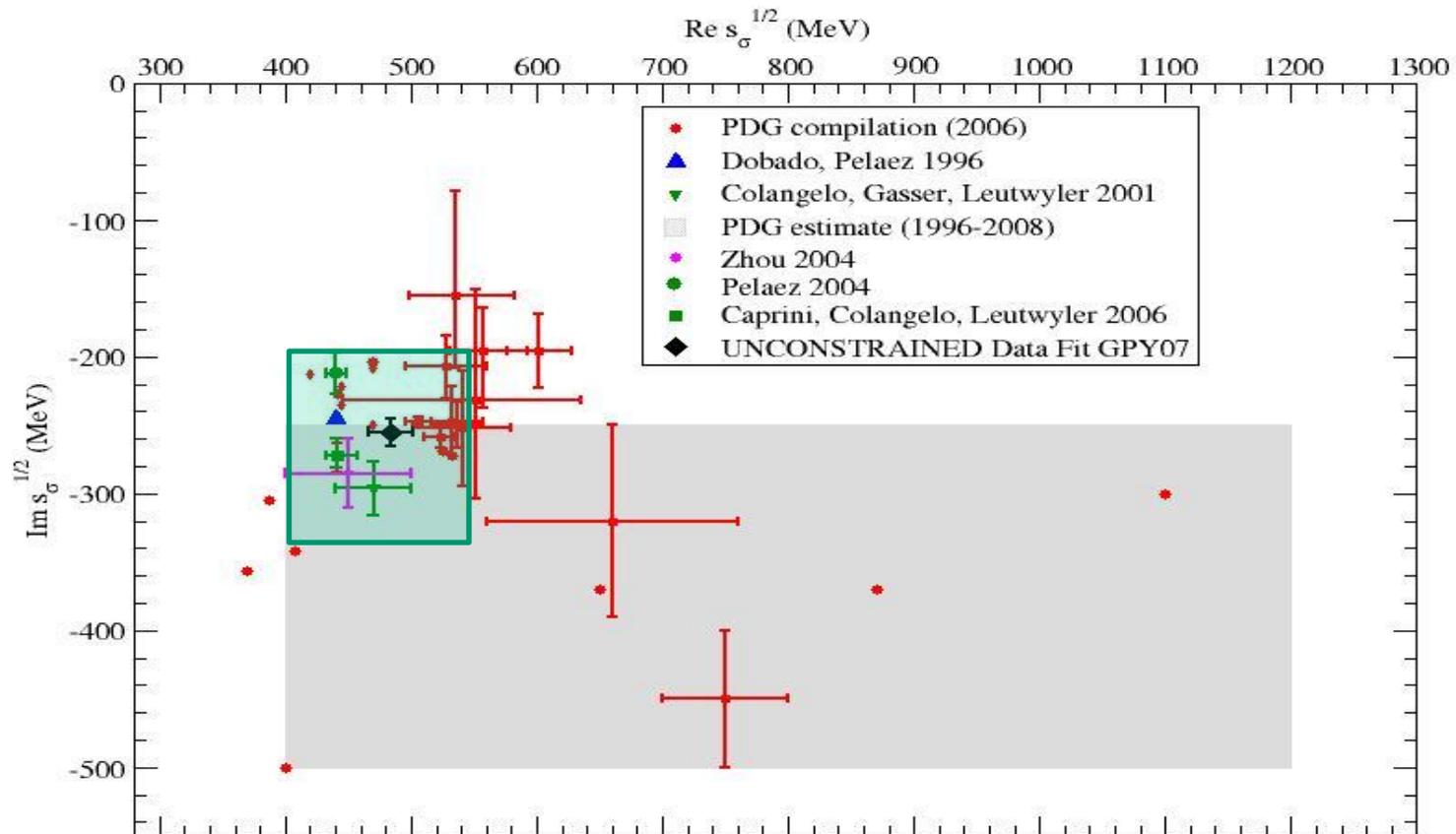
Actually, this happens with the $f_0(500)$ or sigma meson

are doubled due to two flavor components, $n\bar{n}$ and $s\bar{s}$. We do not put the enigmatic σ meson [11–14] on the $q\bar{q}$ trajectory supposing σ is alien to this classification. The broad state

Anisovich-Anisovich-Sarantsev-Phys.Rev.D62.051502-4

Other authors considered that since the $f_0(500)$ had **such a big uncertainty in the PDG it could be ignored.**

Introduction: Regge trajectories



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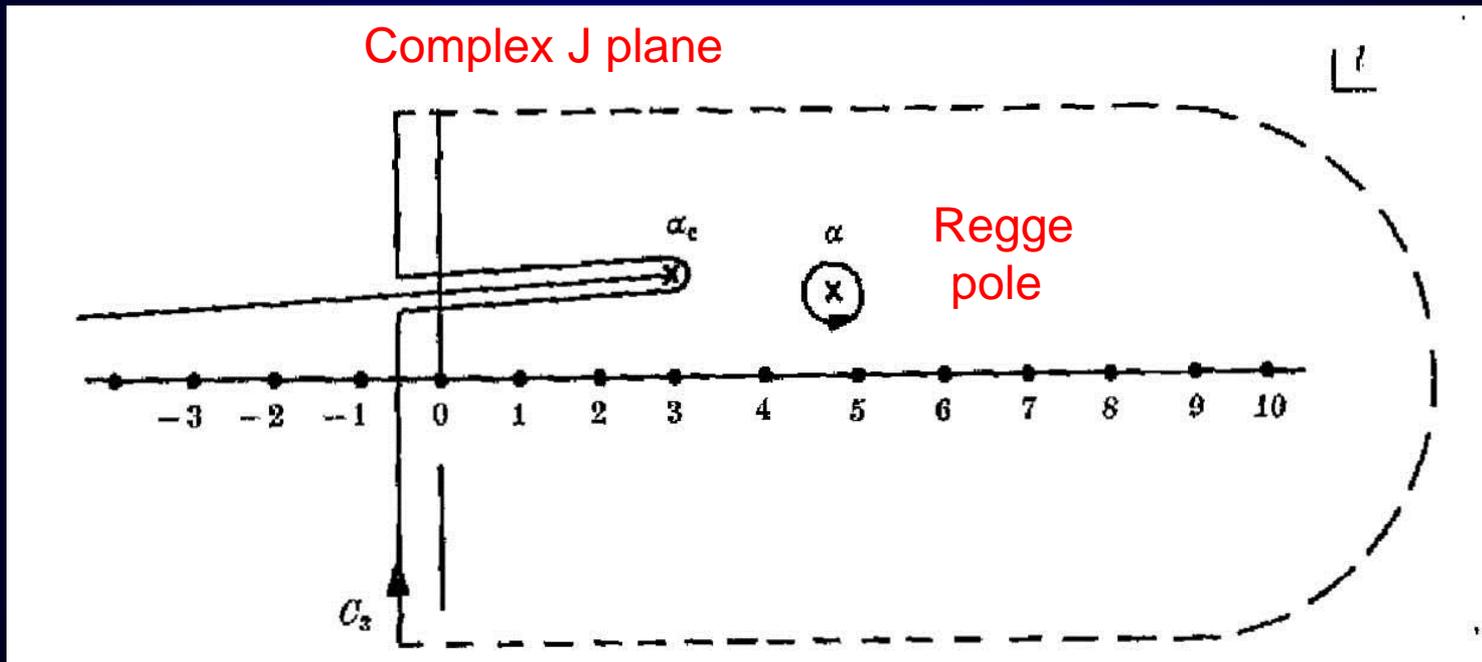
That is no longer an excuse.

One could still think of using the large width as an uncertainty...
But here will study its Regge trajectory as a complex pole,
thus taking into account its width.

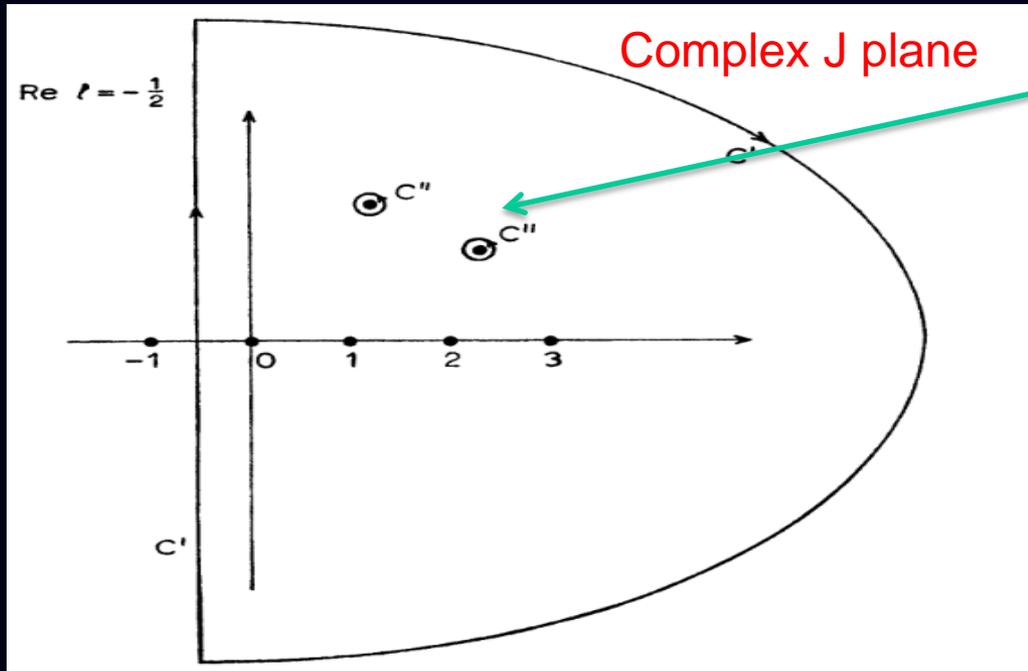
Introduction: Regge Theory

The Regge trajectories can be understood from the analytic extension to the complex angular momentum plane of the partial wave expansion through the Sommerfeld-Watson transform:

$$T(s, t) = \sum_{J=0}^{\infty} (2J+1) f_J(s) P_J(z) \quad \longrightarrow \quad T(s, t) = -\frac{1}{2i} \int_C \frac{(2J+1) f(J, s) P_J(-z)}{\sin \pi J} dJ$$



Introduction: Regge Theory



Regge poles

Position $\alpha(s)$

Residue $\beta(s)$

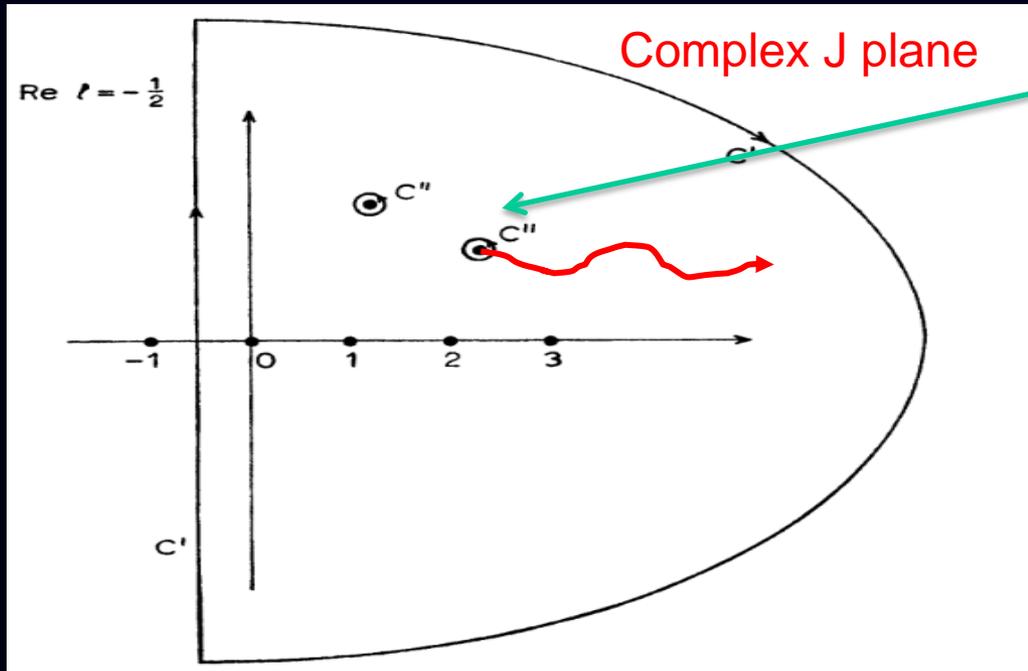
The contribution of a single Regge pole to a partial wave, is shown to be

$$f(J, s) = \hat{f} + \frac{\beta(s)}{J - \alpha(s)}$$

“background” regular function.

Assumption: WE WILL NEGLECT IT in our cases

Introduction: Regge Theory



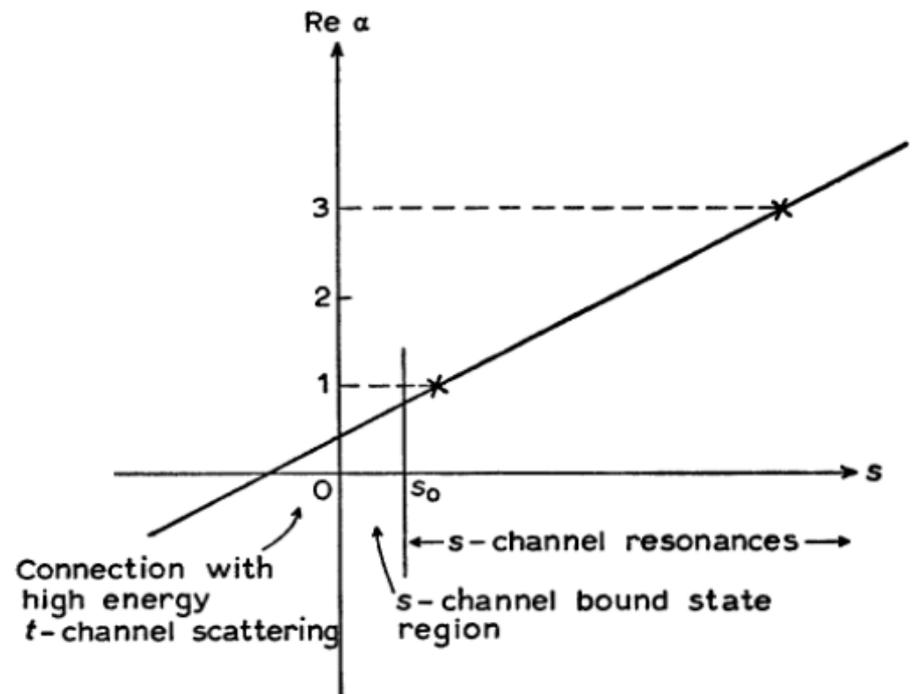
Regge poles

Position $\alpha(s)$

Residue $\beta(s)$

For different s poles move in the complex J plane along Regge Trajectories

Linear trajectories of $q\bar{q}$ mesons are just an example of some specific dynamics



Intriduction: Regge Theory

But other dynamics could lead to different trajectories.

However, trajectories and residues cannot be completely arbitrary due to their analytic properties (Collins, Introduction to Regge Theory)

- **Twice-subtracted dispersion relations**

$$\alpha(s) = A + B(s - s_0) + \frac{(s - s_0)^2}{\pi} \int_{\text{thr.}}^{\infty} \frac{\text{Im}\alpha(s') ds'}{(s' - s)(s' - s_0)^2}$$

$$\gamma(s) = g^2 \exp \left\{ C(s - s_0) + \frac{(s - s_0)^2}{\pi} \int_{\text{thr.}}^{\infty} \frac{\phi_\gamma(s')}{(s' - s)(s' - s_0)} ds' \right\}$$

with

$$\beta(s) = \frac{\hat{s}^{\alpha(s)}}{\Gamma(\alpha(s) + \frac{3}{2})} \gamma(s) = \frac{\text{Im}\alpha(s)}{\rho(s)}$$

Parametrization of pole dominated amplitudes

Chu, Epstein, Kaus, Slansky, Zachariasen, PR175, 2098 (1968).

Moreover, for $\pi\pi$ scattering:

- Unitarity condition on the real axis implies

$$\text{Im } \alpha(s) = \rho(s)\beta(s)$$

$$\rho(s) = \sqrt{1 - 4m_\pi^2/s}$$

- Further properties of $\beta(s)$

threshold behavior

$$\beta(s) = \frac{\hat{s}^{\alpha(s)}}{\Gamma(\alpha(s) + \frac{3}{2})} \gamma(s)$$

$$\hat{s} = \frac{s - 4m^2}{\tilde{s}}$$

suppress poles
of full amplitude

$$(2\alpha + 1)P_\alpha(z_s) \sim \Gamma(\alpha + \frac{3}{2})$$

analytic function:

$\beta(s)$ real on real axis

\Rightarrow phase of $Y(s)$ known

\Rightarrow Omnès-type disp. relation

Parametrization of pole dominated amplitudes

Thus, the trajectory and residue should satisfy a system of integral equations:

$$\begin{aligned}\operatorname{Re}\alpha(s) &= \alpha_0 + \alpha' s + \frac{s}{\pi} PV \int_{4m_\pi^2}^{\infty} ds' \frac{\operatorname{Im}\alpha(s')}{s'(s' - s)}, \\ \operatorname{Im}\alpha(s) &= \rho(s) b_0 \frac{\hat{s}^{\alpha_0 + \alpha' s}}{|\Gamma(\alpha(s) + \frac{3}{2})|} \exp(-\alpha' s [1 - \log(\alpha' \tilde{s})]) \\ &+ \frac{s}{\pi} PV \int_{4m_\pi^2}^{\infty} ds' \frac{\operatorname{Im}\alpha(s') \log \frac{\hat{s}}{\hat{s}'} + \arg \Gamma(\alpha(s') + \frac{3}{2})}{s'(s' - s)}\end{aligned}$$

THESE ARE THE EQUATIONS WE HAVE TO SOLVE

In order to have a consistent trajectory

In the scalar case a slight modification is introduced (Adler zero)

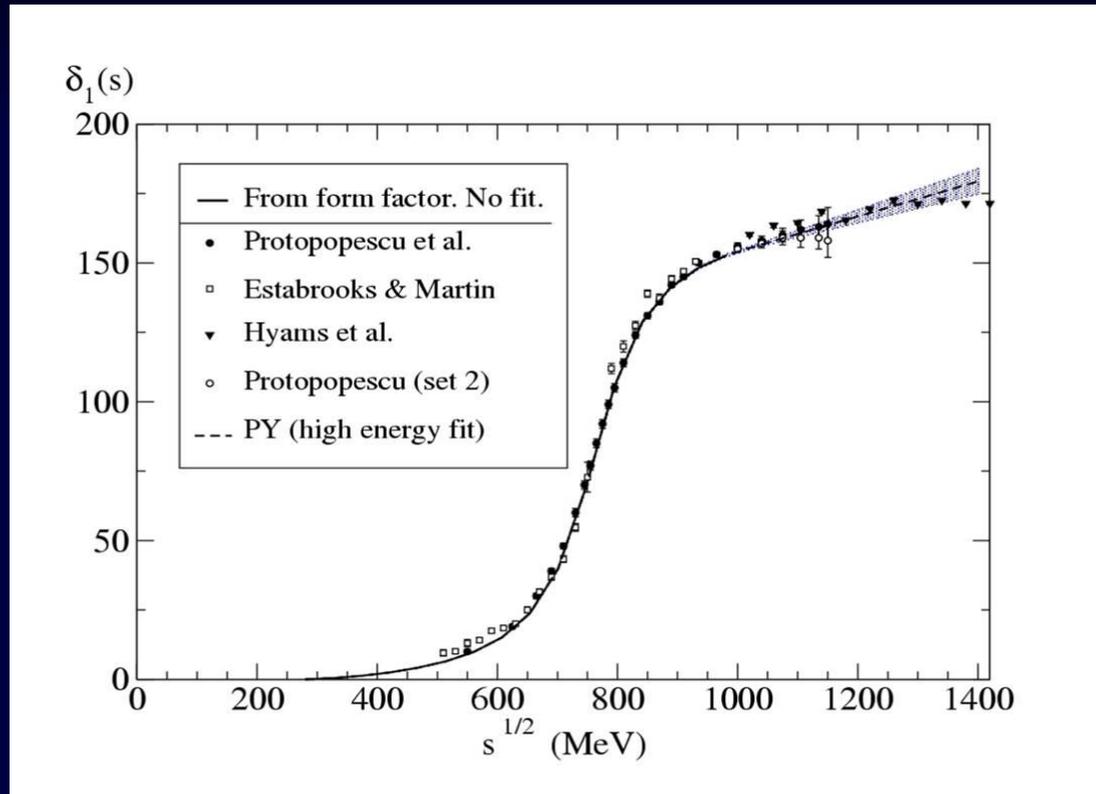
Our Approach

Fix the subtraction constants JUST from the scattering pole

- for a given set of α_0 , α' and b_0 :
 - solve the coupled equations
 - get $\alpha(s)$ and $\beta(s)$ in real axis
 - extend to complex s -plane
 - obtain pole position and residue

$$f^{II}(J, s) = \hat{f} + \frac{\beta(s)}{J - \alpha^{II}(s)}$$

- fit α_0 , α' and b_0 so that the **pole position and residue** coincide with those given by a **dispersive analysis of scattering data**

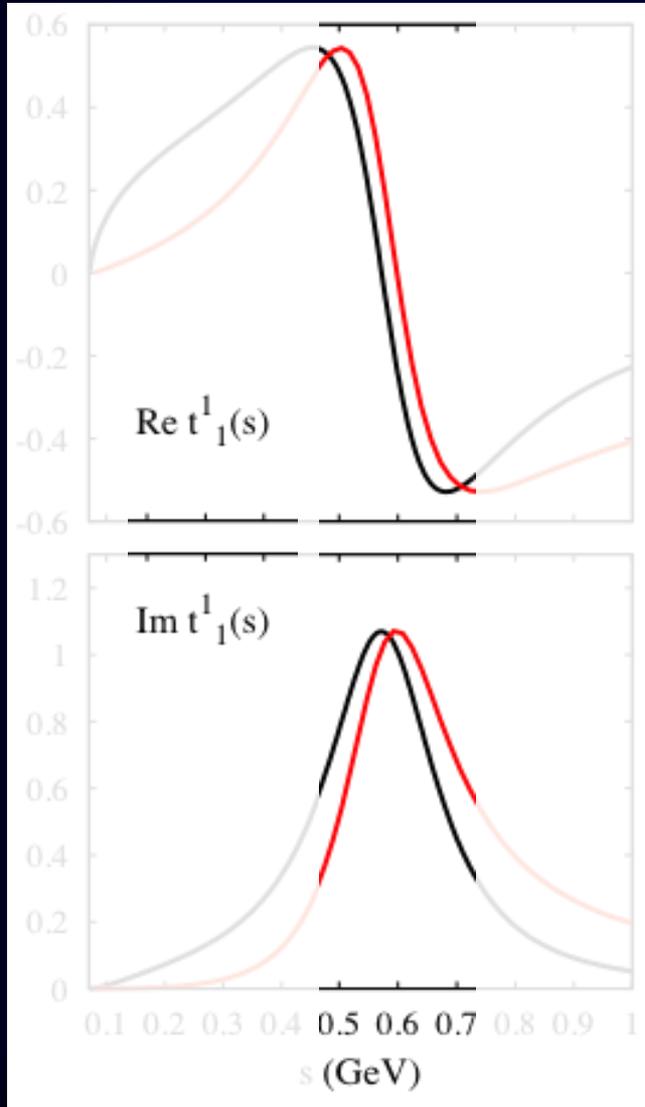


● INPUT for our purposes: **The ρ pole:**

$$\rho_{pole} \approx 763_{-1.5}^{+1.7} - i73.2_{-1.1}^{+1.0} \text{ MeV}$$

$$|g| = 6.01_{-0.07}^{+0.04}$$

Results: ρ case ($l = 1, J = 1$)



We (black) recover a fair representation of the partial wave, in agreement with the GKPY amplitude (red)

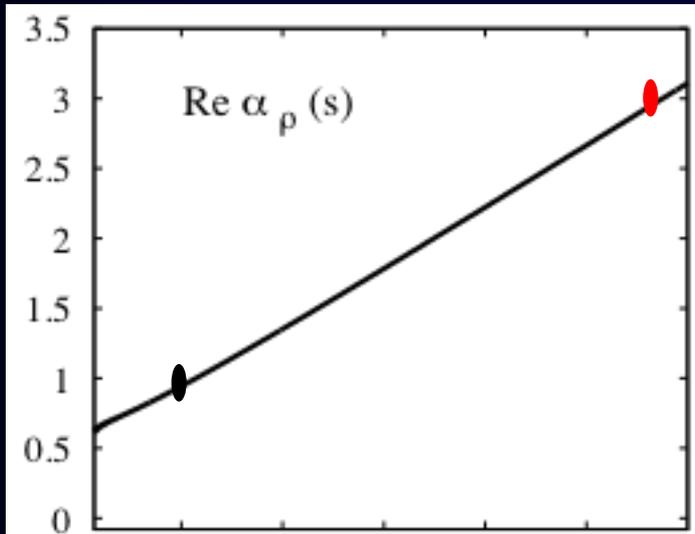
Neglecting the “background” vs. Regge pole gives a 10-15% error.

Particularly in the resonance region

Fair enough to look for the Regge trajectory

Results: ρ case ($l = 1, J = 1$)

We get a prediction for the ρ Regge trajectory, which is almost real



This is a “prediction” for the whole tower of $\rho(770)$ Regge partners:

$\rho(1690)$
 $\rho(2350)$

....

the **LINEAR** behavior
is a **RESULT**

almost LINEAR $\alpha(s) \sim \alpha_0 + \alpha' s$

intercept $\alpha_0 = 0.52$

slope $\alpha' = 0.913 \text{ GeV}^{-2}$

Previous studies:

[1] $\alpha_0 = 0.5$

[1] $\alpha' = 0.83 \text{ GeV}^{-2}$

[2] $\alpha_0 = 0.52 \pm 0.02$

[2] $\alpha' = 0.9 \text{ GeV}^{-2}$

[3] $\alpha_0 = 0.450 \pm 0.005$

[4] $\alpha' = 0.87 \pm 0.06 \text{ GeV}^{-2}$

Remarkably consistent with the literature!!,
(taking into account our approximations)

Results: $f_2(1275)$ and $f_2'(1525)$ cases ($l = 0, J = 2$)

These resonances are almost elastic:

$f_2(1275)$ has BR ($\pi \pi$) = 85% and $f_2'(1525)$ has BR(KK)=90%.

Approximating them with elastic BW-like poles, and solving the integral equations we “predict” again:

almost real and LINEAR $\alpha(s) \sim \alpha_0 + \alpha' s$

For the $f_2(1275)$

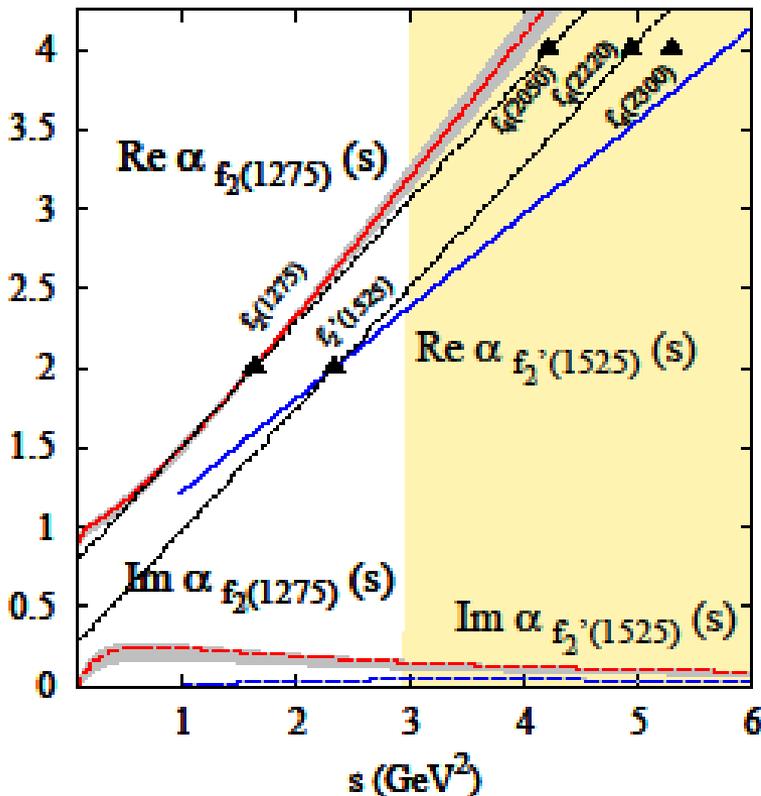
intercept $\alpha_0 = 0.71$

slope $\alpha' = 0.83 \text{ GeV}^{-2}$

For the $f_2'(1525)$

intercept $\alpha_0 = 0.59$

slope $\alpha' = 0.61 \text{ GeV}^{-2}$



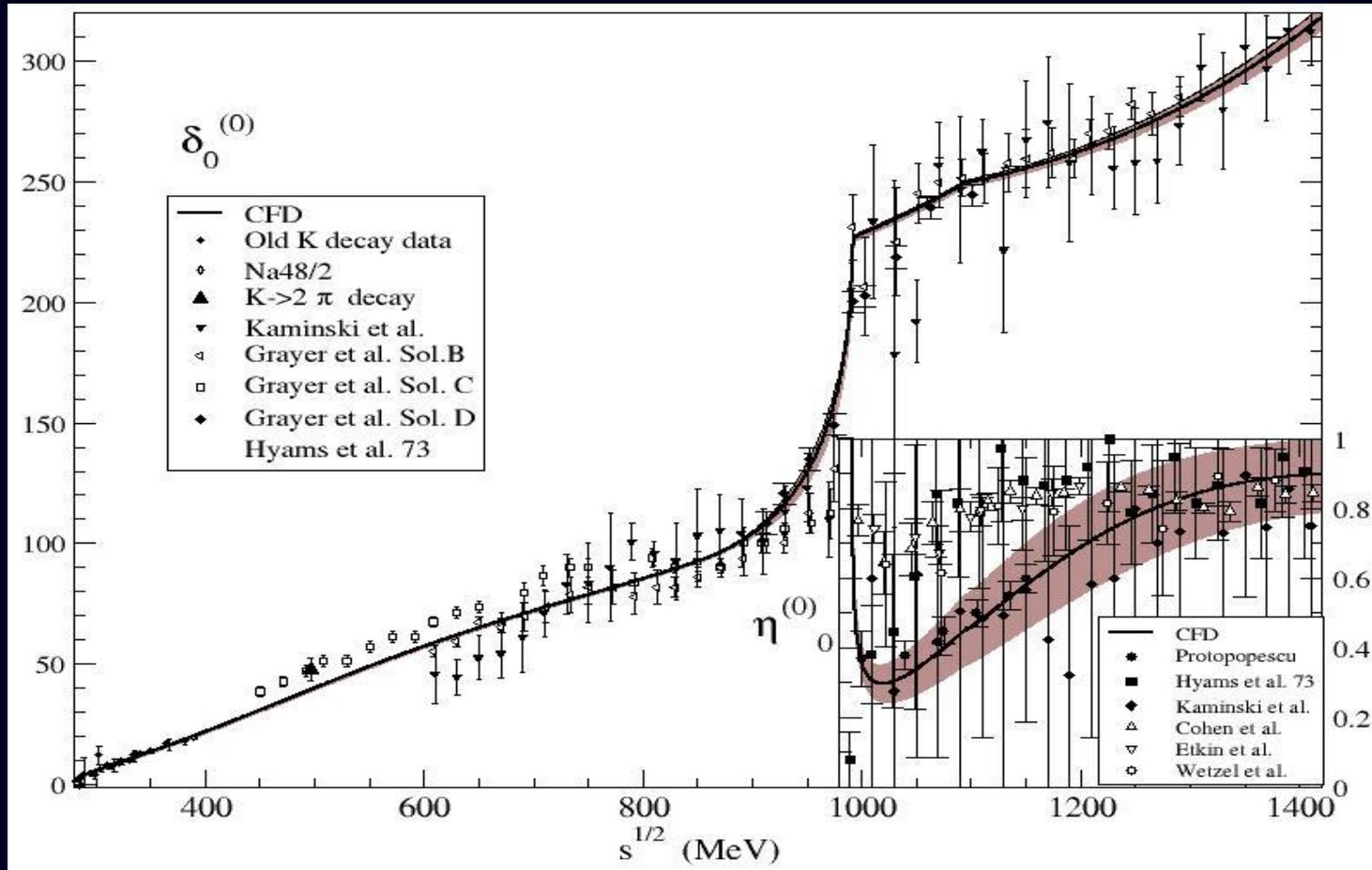
Fair agreement with the literature!!,
(taking into account our approximations)
Remember this is **NOT** a fit!!

This “prediction” for the rho trajectory
Was known since the 70’s, we have just updated it
and obtained new “predictions” for the f_2 and f_2'

So, once we have checked that our approach
Predicts the established Regge trajectories just from the pole
position and residue...

What about the $f_0(500)$?

INPUT: Analytic continuation to the complex plane of a dispersive analysis of data

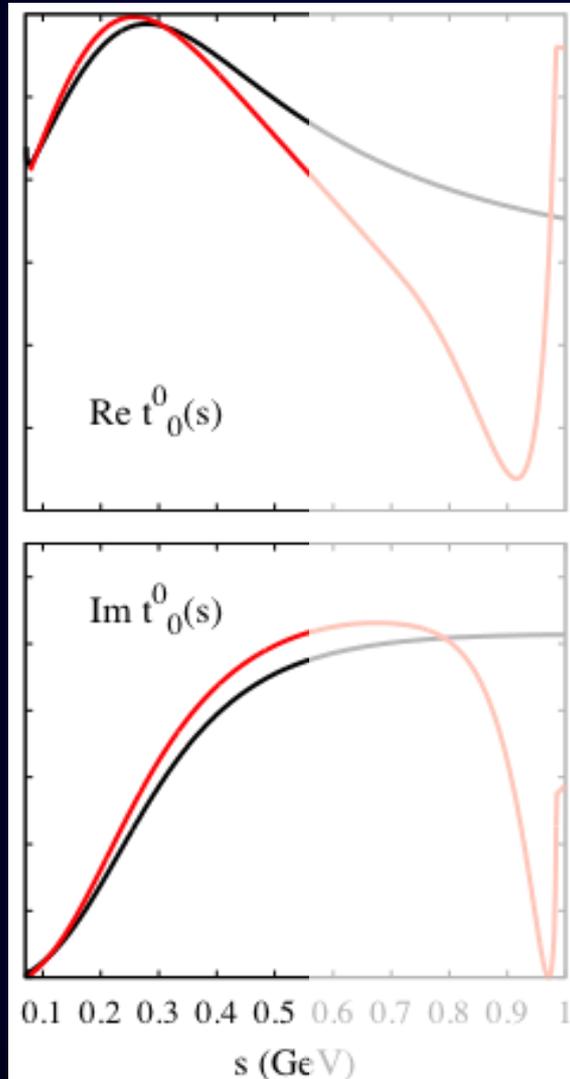


INPUT for our purposes: **The σ pole:**

$$(457_{-15}^{+14}) - i(279_{-7}^{+11}) \text{ MeV}$$

$$|g| = 3.59_{-0.13}^{+0.11} \text{ GeV}$$

Results: σ case ($l = 0, J = 0$)

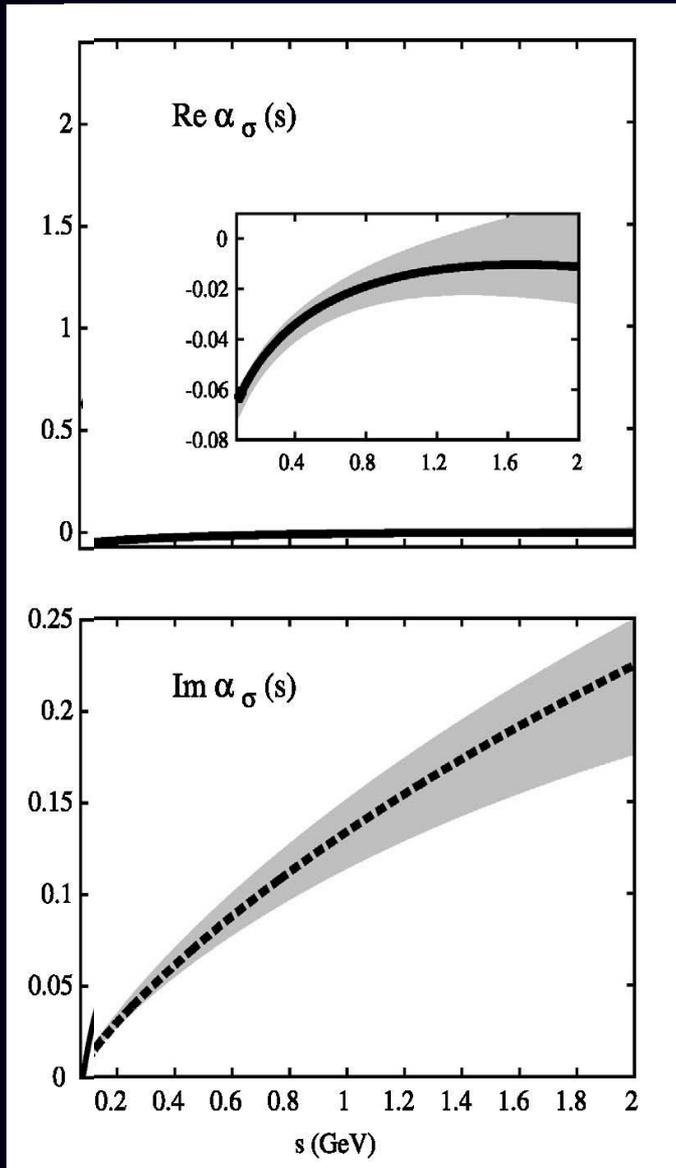


Somewhat better agreement in the resonance region of the Regge pole dominated amplitude with the dispersive amplitude.

So, we apply a similar procedure but now for the $f_0(500)$

Results: σ case ($l = 0, J = 0$)

The prediction for the σ Regge trajectory, is:



- NOT approximately real
- NOT linear

intercept

$$\alpha_\sigma(0) = -0.090^{+0.004}_{-0.012},$$

slope

$$\alpha'_\sigma \simeq 0.002^{+0.050}_{-0.001} \text{ GeV}^{-2}$$

Compare with the rho result...

$$\alpha_0 = 0.52$$

$$\alpha' = 0.913 \text{ GeV}^{-2}$$

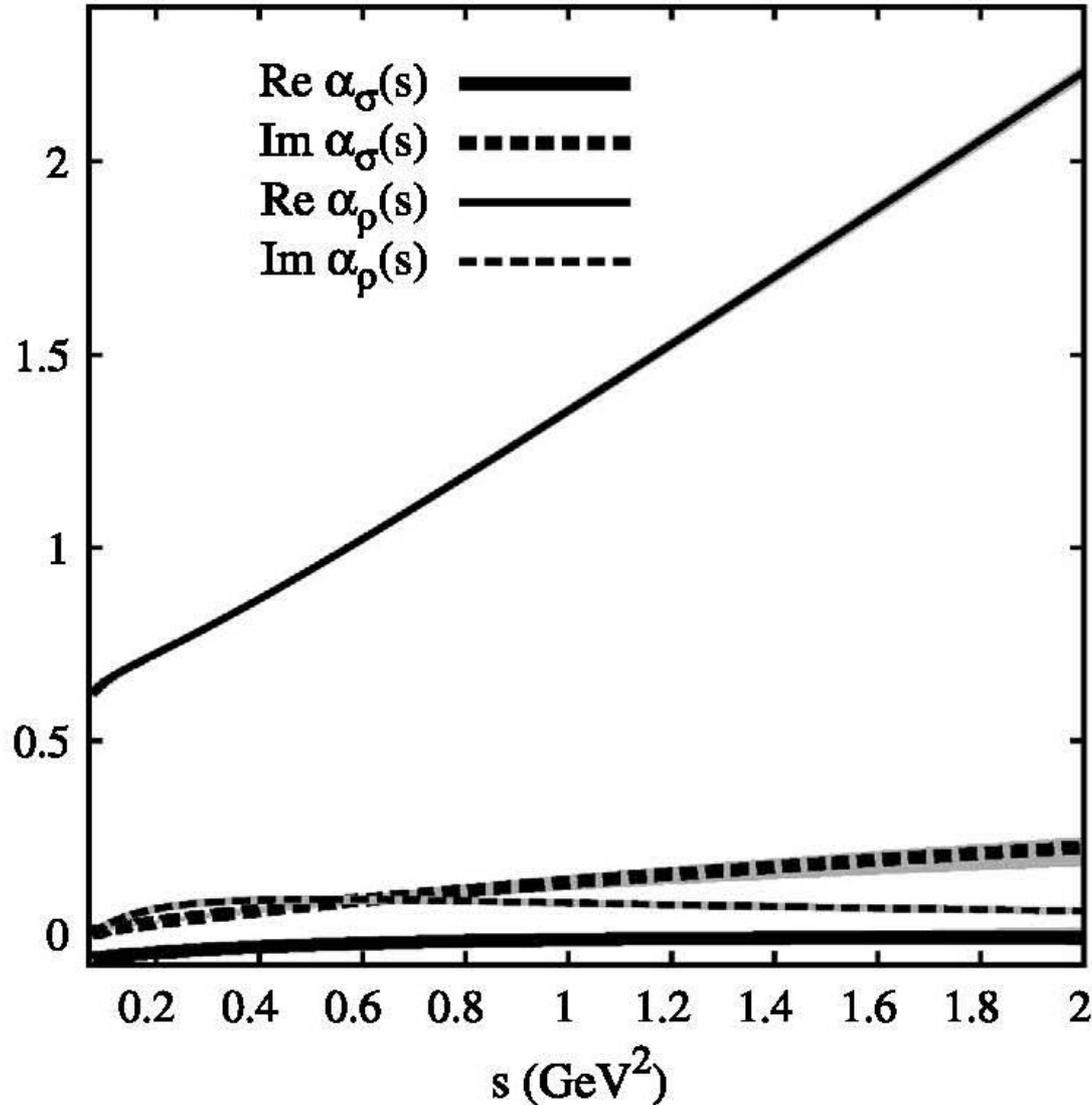
The sigma does **NOT** fit the ordinary meson trajectory

Two orders of magnitude flatter than other hadrons
Typical of meson physics?

$$F_\pi, m_\pi?$$

Results: σ vs. ρ trajectories

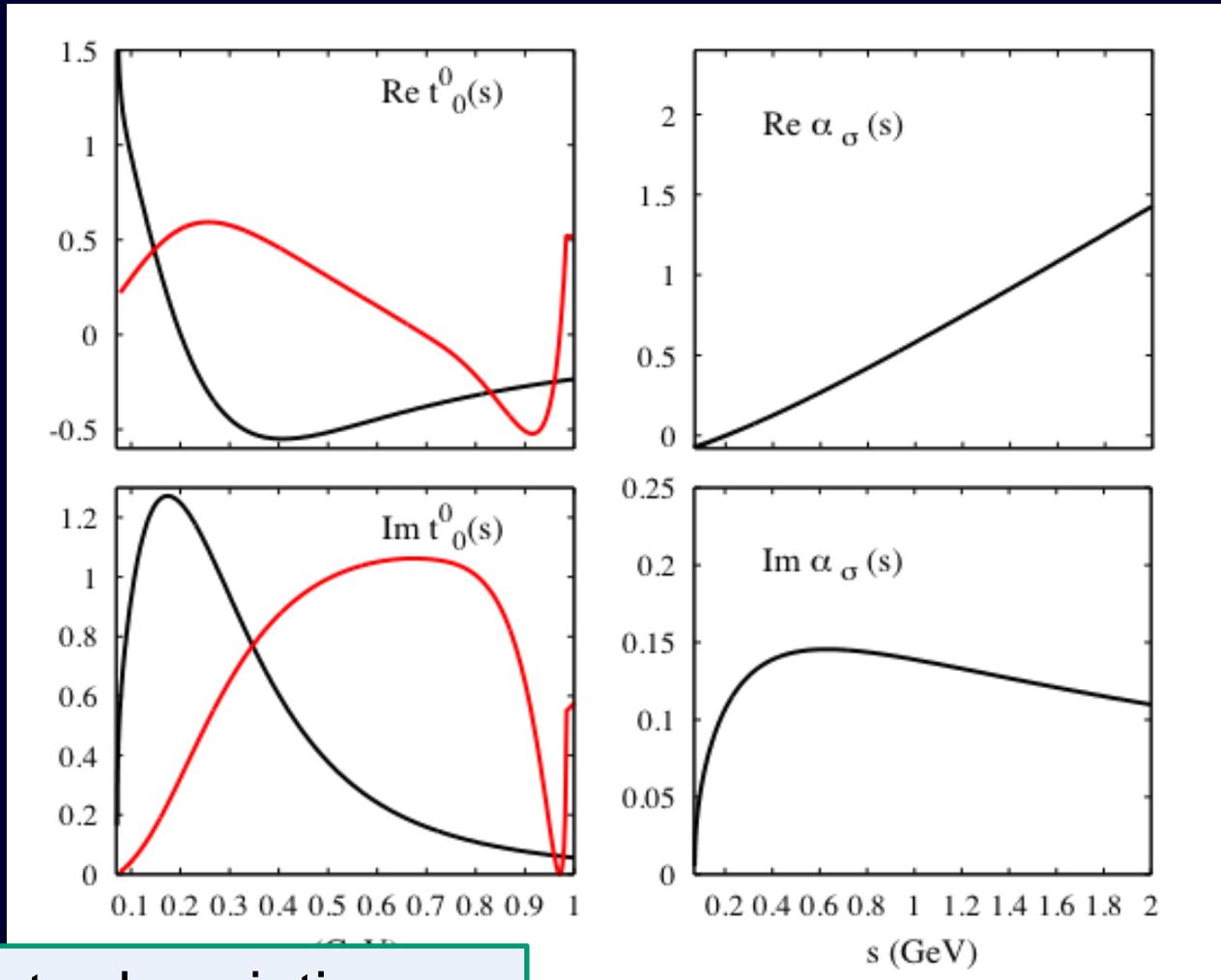
Using the same scale....



No evident
Regge partners
for the $f_0(500)$

Results: σ case ($I = 0, J = 0$)

IF WE INSISTED in fixing the α' to an “ordinary” value $\sim 1 \text{ GeV}^{-2}$



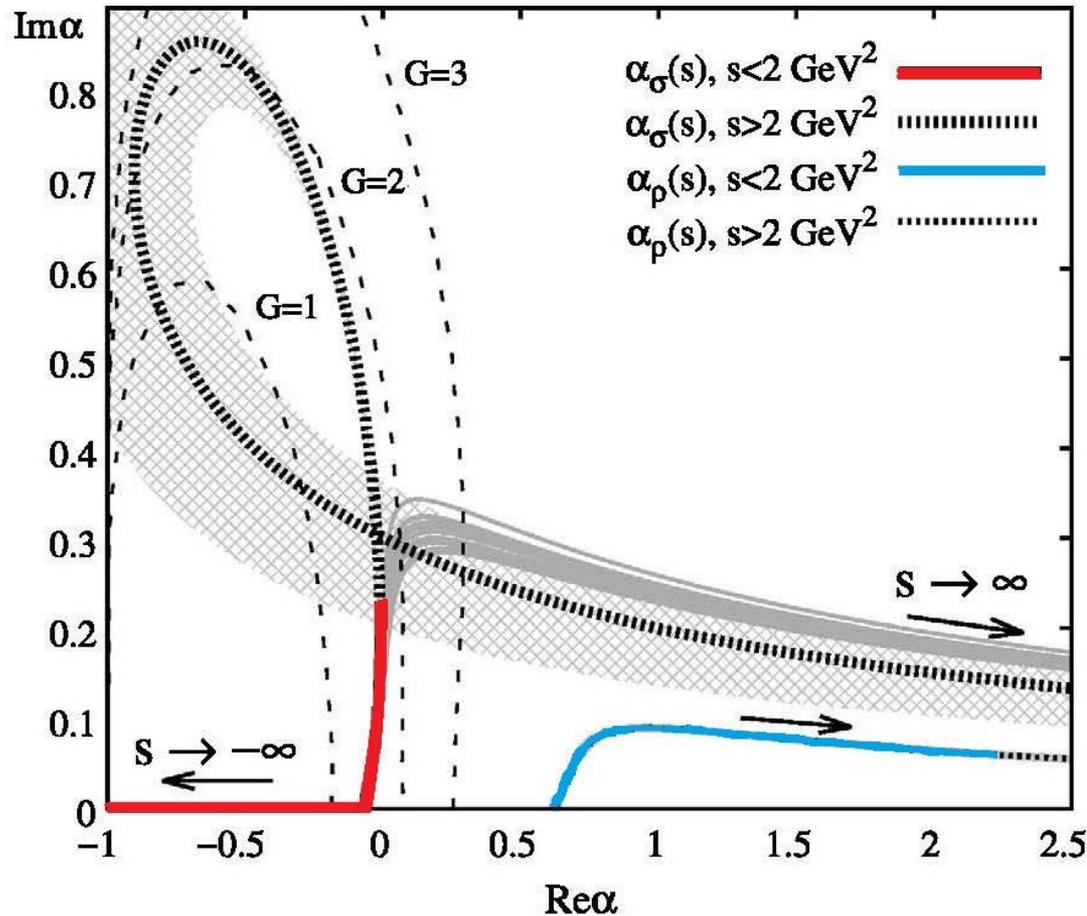
The data description
would be severely spoilt

If not-ordinary...

What then?
Can we identify the dynamics of the trajectory?

Not quite yet... but...

Plotting the trajectories in the complex J plane...



Striking similarity with Yukawa potentials at low energy: $V(r) = -Ga \exp(-r/a)/r$

Our result is mimicked with $a = 0.5 \text{ GeV}^{-1}$ to compare with S-wave $\pi\pi$ scattering length 1.6 GeV^{-1}

“a” rather small !!!

Non-ordinary σ trajectory

The extrapolation of our trajectory also follows a Yukawa but deviates at very high energy

Ordinary ρ trajectory

The extrapolation of our trajectory also follows a Yukawa but deviates at very high energy

Summary

- Analytic constraints on Regge trajectories as integral equations.
- Consistent treatment of the width
- **Fitting JUST the pole position and residue of an isolated resonance,** yields its Regge trajectory parameters
- ρ , f_2 and f_2' trajectories: COME OUT LINEAR, with universal parameters
- σ trajectory: NON-LINEAR.
 - Trajectory slope **two orders of magnitude smaller**
 - No partners.
- If we force the σ trajectory to have universal slope, data description ruined
- At low energies, striking similarities with trajectories of Yukawa potential

Outlook

- Different mass case (κ , K^* ...)
- Meson-nucleon (Delta, etc...)
- Similar non-linear behavior for other non-ordinary states?
- Coupled channels
- Microscopic models, relation to compositeness, etc...

**Happy Birthday Eef and
congratulations for these
*first 70 years of very fruitful career***

**Thank you for your advice and
personal example of dedication and
enjoyment when doing Physics**